ELSEVIER



Journal of Alloys and Compounds



journal homepage: www.elsevier.com/locate/jallcom

Effects of Fe substitution on structural and magnetic properties of the $Nd_2Co_{7-x}Fe_x$ compounds

YuQi Yang, G.H. Rao*, T. Wang, J.B. Li, J. Luo, G.Y. Liu, X.J. Chen, J.L. Zhao

Beijing National Laboratory for Condensed Matter Physics, Institute of Physics, Chinese Academy of Sciences, Beijing 100190, China

A R T I C L E I N F O

Article history: Received 13 May 2010 Received in revised form 4 July 2010 Accepted 7 July 2010 Available online 15 July 2010

Keywords: Nd₂Co_{7-x}Fe_x Crystal structure Spin-reorientation Magnetic phase diagram

ABSTRACT

The structural and magnetic properties of the Nd₂Co_{7-x}Fe_x compounds related to $R_{m+n}T_{2m+5n}$ series (R = rare earth, T = transition metal) were investigated by means of X-ray diffraction and magnetic measurements. The compounds crystallize in the Ce₂Ni₇-type structure with the space group *P*6₃/*mmc* between *x* = 0.0 and 2.1. For higher Fe content the Nd₂(CoFe)₇ either coexists with the Nd(CoFe)₃ phase (*x* > 2.1) or decomposes to Nd(CoFe)₃ + Nd₂(CoFe)₁₇ (*x* > 2.4). The single phase Nd₂Co_{7-x}Fe_x compounds exhibit two spin-reorientation transitions (SRT) for *x* < 1.5. The substitution of Fe for Co elevates both the spin-reorientation transition temperatures for *x* < 1.5 and finally leads to a single ferromagnetism to paramagnetism transition for *x* ≥ 1.5. A magnetic phase diagram is constructed based on the magnetic cally aligned samples. The saturation magnetization of the compounds increases gradually first with the Fe content for *x* ≤ 0.20, substantially for 0.20 < *x* < 1.5, and then decreases for *x* > 1.5, which can be elucidated by a depletion effect of d band electrons due to the substitution of Fe for Co and the spin-flipping due to an enhanced d band splitting.

© 2010 Elsevier B.V. All rights reserved.

1. Introduction

The rare earth cobalt compounds with the hexagonal Ce₂Ni₇type structure or the rhombohedral Gd₂Co₇-type structure belong to a structural series $R_{m+n}T_{2m+5n}$ (R = rare earth atoms, T = transition metal atoms), which is composed of m layers of Laves-phase RT₂ slabs and n layers of CaCu₅-type slabs stacking alternately along the *c* direction [1]. The stacking of the slabs is accompanied by a shift of the layers that gives rise to either a hexagonal or a rhombohedral symmetry [2,3]. With high magnetocaloric effect (MCE) of the Laves-phase RCo₂ at low temperature [4–7] and high saturation magnetization of CaCu₅-type phase at high temperature, $R_{m+n}T_{2m+5n}$ compounds might be promising MCE candidates for applications around ambient temperature.

The binary compound Nd₂Co₇ (m = 2, n = 2) is stable in the Nd–Co system and crystallizes in the hexagonal Ce₂Ni₇-type structure with the space group *P*6₃/*mmc*. Andreev et al. [8] and Bartashevich et al. [9] reported that the magnetic structure of Nd₂Co₇ experienced two fluctuations as temperature increases: a spinreorientation (SR) occurred around 225 K and another SR around 290 K, with an alteration of the easy magnetization direction (EMD) from *a*-axis below the first SR transition to the *c*-axis above the second SR transition. However, binary " R_2Fe_7 " does not exist possibly due to much higher mixing enthalpy between R and Fe than that between R and Co [10]. In addition, Co and Fe atoms on the same crystallographic position in a rare earth transition metal intermetallic compound usually exhibits opposite sign of magnetocrystalline anisotropy and different ferromagnetic character (strong vs. weak ferromagnetism). Therefore, an exploration of new members of the $R_{m+n}T_{2m+5n}$ family with Fe element is of significance for enriching our understanding of rare earth transition metal intermetallic compounds. In this paper we report crystal structure and magnetic properties of $Nd_2Co_{7-x}Fe_x$ compounds. Special attention will focus on the solubility limit of Fe and the influence of Fe atom on the structure, magnetic properties and spin-reorientation transitions (SRT) of the compounds.

2. Experimental details

Alloys with the composition Nd₂Co_{7-x}Fe_x (x=0.0-4.0) were prepared by arcmelting metal ingots of Nd, Co and Fe (purity >99.98% for Fe and Co, >99.9% for Nd) in an argon atmosphere. All alloys were remelted at least four times to ensure homogeneity. The alloy ingots were wrapped in Ta foil, sealed into evacuated quartz tubes and annealed one week in vacuum at 1353 K for x=0.0-1.0, at 1303 K for x=1.0-2.0, and at 1263 K for x=2.0-4.0, respectively, then furnace-cooled to room temperature.

The samples were examined by means of X-ray powder diffraction (XRD) and thermo-magnetic analysis. The XRD data were collected on a Rigaku D/max-2500 diffractometer with Cu K α radiation and a graphite monochromator. The XRD data used for structure refinement were collected in a step-scan mode with a sampling

^{*} Corresponding author at: Beijing National Laboratory for Condensed Matter Physics, Institute of Physics, Chinese Academy of Sciences, Beijing 100190, China.

E-mail addresses: yangyuqi5@gmail.com (Y.Q. Yang), ghrao@aphy.iphy.ac.cn (G.H. Rao).

^{0925-8388/\$ -} see front matter © 2010 Elsevier B.V. All rights reserved. doi:10.1016/j.jallcom.2010.07.065



Fig. 1. XRD patterns of $Nd_2Co_{7-x}Fe_x$ ($0 \le x \le 2.5$).

time of 3 s and a sampling step of 0.02° in the 2θ range of $18-120^{\circ}$. Temperature dependence of magnetization of the samples was measured under a field of 1 kOe with a vibrating-sample magnetometer (VSM) from room temperature to above the Curie temperature. The ac magnetic susceptibility below 300 K of the samples was measured by a mutual inductance bridge at a fixed frequency of 320 Hz. Field dependence of magnetization of the samples was measured on a SQUID magnetometer at 5 K in magnetic fields up to 70 kOe, and the saturation magnetization was derived according to the law of approach to saturation. The Curie temperature was determined based on thermogravimetry analysis (SDT Q600 TA Instrument) under a low magnetic field exerted by a permanent magnet. The easy magnetization direction of the samples at room temperature was identified on the basis of the XRD pattern of magnetically pre-aligned specimens, which were prepared by solidifying a mixture of fine alloy particles with epoxy resin at room temperature in a field of about 10 kOe perpendicular to the specimen surface.

3. Results and discussion

3.1. Crystal structure

Portions of the XRD patterns of the $Nd_2Co_{7-x}Fe_x$ are shown in Fig. 1 for phase identification. The sample is single phase and crystallizes in the Ce₂Ni₇-type structure (space group $P6_3/mmc$) for

Lattice parameters *a* and *c*, Curie temperature T_{C} , saturation magnetization M_{S} at 5 K, and the EMD at room temperature for Nd₂Co_{7-x}Fe_x compounds.

x	<i>a</i> (Å)	<i>c</i> (Å)	<i>T</i> _C (K)	EMD	$M_{\rm S}$ ($\mu_{\rm B}/{ m f.u.}$)
0.0	5.0648(2)	24.4652(7)	605.2	c-Axis	13.60
0.2	-	-	622.2	c-Axis	13.82
0.4	5.0702(1)	24.4413(8)	638.8	Cone	14.24
0.6	-	-	652.5	Cone	14.64
0.8	5.0809(2)	24.4233(7)	666.6	Planar	14.98
1.0	5.0898(1)	24.4517(8)	680.3	Planar	15.85
1.2	5.0914(3)	24.4321(11)	689.2	Planar	16.34
1.4	5.0968(2)	24.4301(7)	698.4	Planar	17.06
1.5	5.1010(3)	24.4486(12)	701.4	Planar	17.17
1.8	5.1063(2)	24.4307(11)	714.1	Planar	16.98
2.1	5.1093(3)	24.4238(13)	725.1	Planar	-
2.4	5.1155(2)	24.4308(13)	730.3	Planar	-

x = 0.0-2.1. With the increase of the Fe content, the intensity of a shoulder peak at $2\theta \sim 32.5^{\circ}$ increases for x = 2.1-2.5 (Fig. 1), indicating an increasing amount of the RT₃ phase (also a member of $R_{m+n}T_{2m+5n}$ -type compounds with m = 2 and n = 1). For the sample with x = 2.5, the 2:7 phase disappears completely and decomposes to Nd(CoFe)₃ + Nd₂(CoFe)₁₇. The lattice parameters of the compounds were derived by profile decomposition and listed in Table 1. Fig. 2 shows the composition dependence of lattice constants for the Nd₂(CoFe)₇ phase (x < 2.5), the Nd(CoFe)₃ phase ($x \ge 2.5$) and the Nd₂(CoFe)₁₇ phase (inset). For the 2:7 phase, the lattice parameter *a* increases linearly with the Fe content, whereas the lattice parameter *c* decrease slightly.

The crystal structure of the single phase $Nd_2Co_{7-x}Fe_x$ was refined by the Rietveld technique, using the refinement program Rietan 2000 [11]. Fig. 3 displays the experimental and calculated XRD patterns as well as the difference between them for Nd_2Co_7 and $Nd_2Co_{5.5}Fe_{1.5}$. The atomic parameters for some selected compounds, Nd_2Co_7 , $Nd_2Co_{6.2}Fe_{0.8}$ and $Nd_2Co_{5.5}Fe_{1.5}$, were listed in Table 2.

Because Fe has a larger atomic radius than Co, an expansion of lattice is expected with the increasing substitution of Co by Fe. Based on the refined structural parameters, the thicknesses of the Nd(CoFe)₅ slabs and the Nd(CoFe)₂ slabs along the *c* direction in Nd₂Co_{7–x}Fe_x, i.e. the interplanar distance between atomic layer of the atoms on the 12k sites, can be calculated. With the increase



Fig. 2. The composition dependence of lattice constants of Nd₂Co_{7-x}Fe_x. 2:7 phase for x < 2.5, 1:3 phase and 2:17 phase (inset) for $x \ge 2.5$. The 2:7 phase and 1:3 phase coexist for 2.1 < x < 2.5, and the 1:3 phase and 2:17 phase coexist for $2.5 \le x \le 4$.



Fig. 3. Observed and calculated XRD patterns of the compound Nd₂Co₇ (a) and Nd₂Co_{5.5}Fe_{1.5} (b). In each graph, the lowest curve is the difference between the observed and the calculated intensities, and the vertical bars at the bottom indicate Bragg reflection positions.

of the Fe content, the thickness of the Nd(CoFe)₅ slabs increases, whereas the thickness of the Nd(CoFe)₂ slabs decreases, leading to the observed slight reduction of the *c*-axis in Nd₂Co_{7-x}Fe_x. Therefore, it seems that the substitution of Fe for Co results in an expansion of the Nd(CoFe)₅ slabs along all three directions as in Nd(Co_{0.95}Fe_{0.05})₅ and Y(Co_{1-x}Fe_x)₅ compounds [12–14] and causes a shrinkage of the Nd(CoFe)₂ slabs. The shrinking of the Nd(CoFe)₂ slabs coincides with the fact that the NdFe₂ compound was stabilized under a high pressure [15].

For the CaCu₅-type structure, it was reported that the large Fe atoms occupied preferentially the 3g site in Nd(Co_{0.95}Fe_{0.05})₅ and Y(Co_{1-x}Fe_x)₅ [12–14]. The 3g sites in the CaCu₅-type structure form a kagomè lattice and correspond to the 6 h and the 12k sites in the Nd₂Co₇ structure. It is likely that the Fe atoms also preferentially reside on the kagomè lattice in Nd₂Co_{7-x}Fe_x, i.e. the 6 h and the 12K sites. Neutron diffraction experiments are desirable to determine

Table 2 Rietveld refinement results of the crystal structure of $Nd_2Co_{7-x}Fe_x$, space group $P6_3/mmc$. Co1 at 2a (0.0.0).

$Nd_2Co_{7-x}Fe_x$	<i>x</i> = 0.0	<i>x</i> = 0.8	<i>x</i> = 1.5		
Lattice constants					
a (Å)	5.0648(2)	5.0809(2)	5.1010(3)		
<i>c</i> (Å)	24.4652(7)	24.4233(7)	24.4486(12)		
Atomic positions					
Nd1: 4f(1/3,2/3,z), z	0.0268(1)	0.0269(2)	0.0268(2)		
Nd2: 4f(1/3,2/3,z), z	0.1721(2)	0.1717(1)	0.1719(2)		
Co2: 4e(0,0,z), z	0.1736(4)	0.1708(3)	0.1724(5)		
Co3: 4f(1/3,2/3,z), z	0.8338(3)	0.8320(3)	0.8334(5)		
Co4: 6h (<i>x</i> ,2 <i>x</i> ,1/4), <i>x</i>	0.8309(8)	0.8358(8)	0.8309(14)		
Co5:12k (<i>x</i> ,2 <i>x</i> , <i>z</i>), <i>x</i>	0.8370(6)	0.8393(5)	0.8404(9)		
Z	0.0873(2)	0.0871(2)	0.0864(2)		
R factors:					
R _{wp} (%)	8.81	9.39	8.97		
R _p (%)	6.84	7.24	6.96		
R _e (%)	5.22	4.86	6.97		

the distribution of the Fe atoms over the five crystallographically inequivalent positions.

3.2. Magnetic structure and magnetic properties

The parent compound Nd₂Co₇ experiences two spinreorientation transitions as the temperature increases: one is at T_{SR1} = 218 K and the other at T_{SR2} = 261 K as shown in Fig. 4. This observation is consistent with the report of Andreev et al. [8] and Bartashevich et al. [9], showing that the easy magnetization direction of Nd₂Co₇ single crystal was along the *a*-axis below 225 K and along the *c*-axis above 290 K. In addition, the magnetization curve of Nd₂Co₇ mimics that of NdCo₅, implying a predominant contribution from the 1:5 slabs. The two spin-reorientation transitions are observed in the Fe substituted compounds Nd₂Co_{7-x}Fe_x for *x* < 1.5. However, the plateau defined as $\Delta T = T_{SR2} - T_{SR1}$ is narrowed with the increase of the Fe content and finally disappears for *x* ≥ 1.5.

Fig. 5 shows the representative XRD patterns of the $Nd_2Co_{7-x}Fe_x$ fine powders magnetically aligned at room temperature in a field of ~ 10 kOe. For Nd₂Co₇ and samples with $x \le 0.2$, the (001) reflections are strongly enhanced, indicating an easy magnetization direction (EMD) along the *c*-axis above T_{SR2} (>300 K). For the sample with $x \ge 0.8$, the (hh0) reflections are strongly enhanced, implying that the EMD of the samples is planar at room temperature. For the samples with $0.4 \le x \le 0.6$, the XRD patterns of the magnetically aligned powders are saliently different: both (001) and (hh0) reflections are weakened while the (h01) reflections are enhanced, implying the EMD of the sample at room temperature deviates from *c*-axis and *ab*-plane. Ilyn et al. have reported the anisotropy constants between 210K and 270K of Nd₂Co₇ that well reproduced the observed temperatures of the spin-reorientation transition [16]. From the reported anisotropy constants it is easy to derive that the EMD of Nd_2Co_7 alters rapidly from the *ab*-plane at T_{SR1} to the *c*-axis at T_{SR2} via a cone anisotropy. Considering $T_{SR1} < 300 \text{ K} < T_{SR2}$ for the samples with $0.4 \le x \le 0.6$



Fig. 4. Temperature dependence of ac susceptibility (a) and magnetization (b) of $Nd_2Co_{7-x}Fe_x$ compounds, the curve of $NdCo_5$ was measured and is presented for comparison. The arrow on x = 2.1 curve marks the Curie temperature of $Nd(CoFe)_3$ phase.

(Fig. 4), it is reasonable to infer that these compounds have a cone anisotropy at room temperature.

For both the Nd₂Co₇ and NdCo₅, the planar anisotropy of the Nd sublattice dominates at low temperature, while the *c*-axis anisotropy of the Co sublattice prevails at high temperature. The Stevens factors for the electronic configurations of Co and Fe have different signs in $R_{m+n}T_{2m+5n}$ series compounds [17]. Paoluzi et al. [18] also found that the overall anisotropy of Fe is not only 2–3 times larger than but also always opposite to that of Co in Y₂(CoFe)₇ compounds. In Y(CoFe)₅ structure, Franse et al. [19] reported that the anisotropy constant K₁s were 5.7×10^6 and -1.7×10^6 J/m³ for Co in 2c and 3g sites, respectively, while those for Fe were almost one order of magnitude larger, i.e. -51.3×10^6 and 11.0×10^6 J/m³,



Fig. 5. The XRD patterns of magnetically pre-orientated powder sample of $Nd_2Co_{7-x}Fe_x$. The lowest curve is the XRD pattern of the un-orientated powder sample with x = 0.0.



Fig. 6. Thermogravimetry traces under a low magnetic field for the Nd₂Co_{7-x}Fe_x.

respectively. The sign of the anisotropy of Fe in 2c site is opposite to and about five times larger than that in 3g site, which implies that a small amount of Fe in the 2c sites may offset the effect of the preferential substitution of Fe for Co in 3g sites. The 4e and 4f sites in the 2:7 structure, correspond to 2c sites in the 1:5 structure while the 6 h and 12k sites in the 2:7 structure to the 3g site in the 1:5 structure. Similarly, a small amount of Fe substitution for Co in the 4e and 4f sites could weaken the *c*-axis anisotropy of the transition metal sublattice in Nd₂Co_{7-x}Fe_x substantially, giving rise to the enhancement of the planar anisotropy and the increase of the spin-reorientation temperature. For x > 1.5, the planar anisotropy dominates over the whole temperature range investigated, and no spin-reorientation transition takes place.

The spin-reorientation transition temperature and the Curie temperature of the compounds $Nd_2Co_{7-x}Fe_x$ can be readily determined by thermogravimetry analysis under a low magnetic field (Fig. 6) and ac susceptibility measurements (Fig. 4). The spinreorientation transition and order-disorder magnetic transition $(T_{\rm C})$ give rise to apparent weight-losses of a small piece of sample on the thermogravimetry trace under a low magnetic field exerted by a permanent magnet placed outside the furnace and on top of the crucible position. Based on these experiments and the XRD patterns of the magnetically aligned samples, the magnetic phase diagram is constructed as shown in Fig. 7 for the $Nd_2Co_{7-x}Fe_x$ compounds. It should be noted that the $T_{\rm C}$ derived from magnetization curves measured in 1 kOe (Fig. 4b) is about 15-20 K higher than that determined from thermogravimetry, probably due to temperature calibration of different apparatus and to the much higher magnetic field in the magnetization measurements. Fig. 7 indicates the spinreorientation temperatures determined by thermogravimetry and ac susceptibility, both were performed under a low magnetic field, are coincident with each other.

The field dependence of magnetization of Nd₂Co_{7-x}Fe_x compounds at 5 K is shown in Fig. 8, which was measured on fine particles by the SQUID magnetometer up to 70 kOe. All samples exhibit ferromagnetic character and tend to approach saturation in a field higher than 20 kOe. The saturation magnetization M_S are listed in Table 1 and displayed in Fig. 9. The M_S increases first gradually with the Fe content for $x \le 0.20$, substantially for 0.20 < x < 1.5, and then decreases for x > 1.5. Assuming the theoretical trivalent ion moment ($g_J \mu_B = 3.27 \mu_B$) for Nd and a parallel alignment of the Nd moment and the Co moment, it is derived that the atomic moment of Co is $1.01 \mu_B$ in Nd₂Co₇, which is close to the values in Y₂Co₇ ($\mu_{Co} = 1.06 \mu_B$ [20]) and in La₂Co₇ ($\mu_{Co} = 1.06 \mu_B$ [21,22]). Such a



Fig. 7. Magnetic phase diagram of $Nd_2 Co_{7-x}Fe_x$ compounds. The transition temperatures were derived from thermogravimetry and ac susceptibility measurements under a low field.



Fig. 8. The field dependence of magnetization of $Nd_2Co_{7-x}Fe_x$ compounds at 5 K.



Fig. 9. Composition dependence of saturation moment of $Nd_2Co_{7-x}Fe_x$ compounds at 5 K. Inset: composition dependence of the magnetization at 5 K of the magnetically pre-aligned samples with x = 0.6, 0.8 and 1.5.

low atomic moment for Co implies the R_2Co_7 compounds (R = Nd, Y, La) are weak ferromagnetic within the scenario of the magnetic valence model [23–25]. Interestingly, for x < 0.20 the increase of the M_S seems to result essentially from the depletion of spindown electrons due to the substitution of Fe for Co $(dM_S/dx = 1)$. For the sample with x = 1.5, $\mu_{C0} = 1.52 \,\mu_{B}$, which is close to the atomic moment of Co (μ_{Co} = 1.6 μ_{B}) with strong ferromagnetism character [23–25]. Since the atomic moment of transition metal may be underestimated if the crystal-field effect on Nd moment is taken into account, the large atomic moment of the transition metal implies that the compound becomes strong ferromagnetism as the substitution of Fe for Co increases and consequently the reduction of the M_S for x > 1.5 can be attributed essentially to a depletion of the spin-up electrons, i.e. $dM_S/dx = -1$ as shown in Fig. 9. However, for the samples with 0.20 < x < 1.5, the substantial increase of the M_S with the Fe content $(dM_S/dx \gg 1 \text{ and is different from Fe-Co } [26]$ binary alloys) could probably result from both the depletion of d electrons and a spin-flipping of electrons from spin-down to spinup subband. The spin-flipping can occur when the splitting of the d band is increased due to the enhanced coupling between transition metal atoms, which is corroborated by the monotonous increase of the Curie temperature with the Fe content (Fig. 7).

In Fig. 9 the field dependence of magnetization at 5 K is shown in the inset for some samples perpendicularly aligned in a field of 10 kOe at room temperature. For the samples with x = 0.8 and 1.5, the EMD is unchanged upon cooling from room temperature to 5 K, and a field induced spin-reorientation or the first-order magnetization process (FOMP) occurs at 5 K when the applied magnetic field is perpendicular to the alignment direction, which may result in appreciable magnetocaloric effects. In contrast, the EMD of the sample with x = 0.6 changes to planar one at low temperature, so only a small upturn is observed when the applied magnetic field is perpendicular to the alignment direction. Since the FOMP phenomenon in unaligned bulk samples is very weak, the magnetocaloric effects of the bulk samples have not been further pursued.

4. Conclusions

In summary, single phase compounds $Nd_2Co_{7-x}Fe_x$ have been prepared with x < 2.1. The compounds crystallized in the Ce_2Ni_7 type structure with a space group P6₃/mmc. For x > 2.1 the $Nd_2(CoFe)_7$ either coexists with the Nd(CoFe)_3 phase or decomposes to Nd(CoFe)_3 + Nd_2(CoFe)_{17} (x > 2.4). The lattice parameter aof the Nd_2Co_{7-x}Fe_x increases with the Fe content. The thickness of the Nd(CoFe)_5 slabs along the *c*-axis expands while that of the Nd(CoFe)_2 sables shrinks, which leads to a slight decrease of the *c*-axis with the Fe content.

For x < 1.5, the Nd₂Co_{7-x}Fe_x compounds experience two spinreorientation transitions. The substitution of Fe for Co elevates both the spin-reorientation transition temperatures and finally leads to a single ferromagnetism to paramagnetism transition for $x \ge 1.5$. Based on magnetic measurements, thermogravimetry analysis and XRD on magnetically aligned samples, a magnetic phase diagram is constructed for the Nd₂Co_{7-x}Fe_x. The saturation magnetization M_S increases gradually first with the Fe content for $x \le 0.20$, substantially for 0.20 < x < 1.5, and then decreases for x > 1.5. The composition dependence of the M_S can be elucidated by a depletion effect of d band electrons due to the substitution of Fe for Co and the spin-flipping due to an enhanced d band splitting.

Acknowledgements

This work is supported by the National Natural Science Foundation of China (Grants No. 50631040) and the National Basic Research Program of China (Grant No. 2006CB601101, 2006CB605101).

- [11] F. Izumi, T. Ikeda, Mater. Sci. Forum 321-323 (2000) 198.
- [12] J.M. Alameda, D. Givord, C. Jeandey, H.S. Li, Q. Lu, J.L. Oddou, J. Phys. 46 (1985) 1581.
- [13] J. Déportes, D. Givord, J. Schweizer, F. Tasset, IEEE Trans. Magn. 12 (1976) 1000.
- [14] J. Laforest, J.S. Shah, IEEE Trans. Magn. 9 (1973) 217.
- [15] C.C. Tang, W.S. Zhan, Y.X. Li, D.F. Chen, J. Du, G.H. Wu, Z.W. Gan, S.R. Qi, J. Phys. D: Appl. Phys. 31 (1998) 2426.
- [1] E. Parthé, B. Chabot, in: K.A. Sschneidner Jr., L. Eyring (Eds.), Handbook on the [16] Physics and Chemistry of Rare Earths, vol. 7, North-Holland, Amsterdam, 1984, mento, California, May 4-8, 2009.
- Y. Khan, Acta Crystallogr. B 30 (1974) 1533. [2]

References

p. 186.

- [3] K.H.J. Buschow, R.P. van Stapele, J. Appl. Phys. 41 (1970) 4066.
- [4] V.K. Pecharsky, K.A. Gschneidner Jr., J. Magn. Magn. Mater. 200 (1999) 44. [5] P.J. von Ranke, V.K. Pecharsky, K.A. Gschneidner Jr., Phys. Rev. B 58 (1998)
- 12110. [6] P.J. von Ranke, N.A. de Oliveira, V.S.R. de Sousa, D.C. Garcia, I.G. de Oliveira, A.
- Magnus, G. Carvalho, S. Gama, J. Magn. Magn. Mater. 313 (2007) 176. [7] J.C.P. Campoy, E.J.R. Plaza, A.A. Coelho, S. Gama, Phys. Rev. B 74 (2006) 134410.
- [8] A.V. Andreev, M.I. Bartashevich, A.V. Deryagin, S.M. Zadvorkin, Ye. N. Tarasov, Phys. Met. Metallogr. 65 (1988) 61.
- M.I. Bartashevish, T. Goto, M. Yamaguchi, J. Magn. Magn. Mater. 111 (1992) 83.
- [10] G.H. Rao, S. Wu, X.H. Yan, Y.L. Zhang, W.H. Tang, J.K. Liang, J. Alloys Compd. 202 (1993) 101.

- M. Ilyn, A.V. Andreev, E.A. Tereshina, M.I. Bartashevich, V. Zhukova, A. Zhukov, J. Gonzalez, in: The 2009 International Magnetics (Intermag) Conference Sacra-
- [17] M.I. Bartashevich, T. Goto, K. Koui, Phys. B: Condens. Matter 292 (2000) 9.
- [18] A. Paoluzi, L. Pareti, M. Solzi, F. Albertini, J. Magn. Magn. Mater. 132 (1994) 185.
- [19] J.J.M. Franse, N.P. Thuy, N.M. Hong, J. Magn. Magn. Mater. 72 (1988) 361.
- [20] E. Koen, J. Schweizer, F. Tasset, Phys. Rev. 186 (1969) 479.
- [21] W.A.J.J. Velge, K.H.J. Buschow, J. Appl. Phys. 39 (1968) 1717.
- [22] K.H.J. Buschow, J. Less Common Met. 33 (1973) 311.
- [23] A.P. Malozemoff, A.R. Williams, V.L. Moruzzi, Phys. Rev. B 29 (1984) 1620.
- [24] A.R. Williams, V.L. Moruzzi, A.P. Malozemoff, K. Terakura, IEEE Trans. Magn. 19 (1983) 1983.
- [25] G.H. Rao, J. Magn. Magn. Mater. 139 (1995) 204.
- [26] A. Díaz-Ortiz, R. Drautz, M. Fähnle, H. Dosch, J.M. Sanchez, Phys. Rev. B 73 (2006) 224208.